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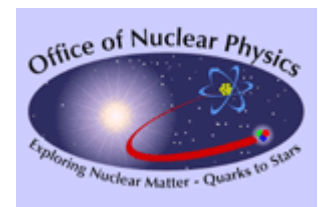
Towards a Faddeev Description of (d,p) Reactions: Separabilization of Optical Potentials

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TORUS collaboration

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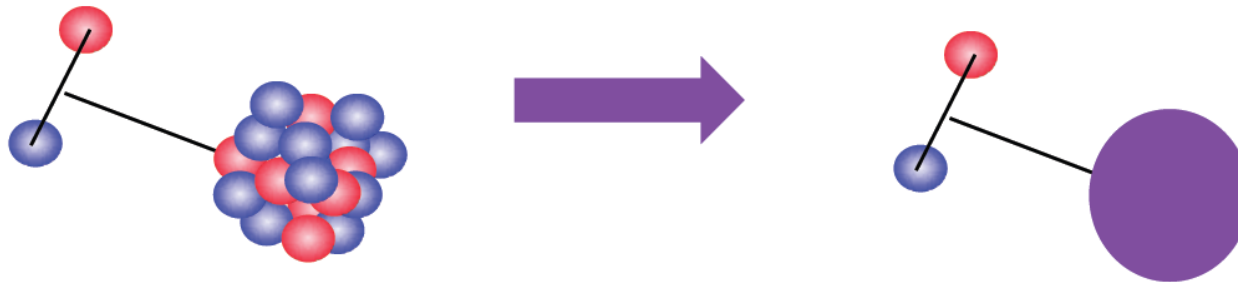


Reactions of Interest:



Ansatz:

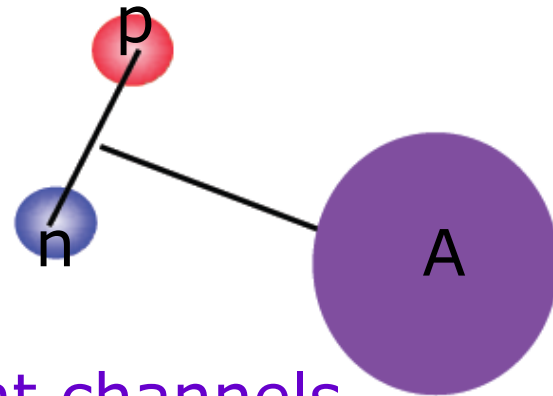
Reduce many-body to few-body problem



- Isolate important degrees of freedom in a reaction
 - Keeping track of all relevant channels
- Connect back to the many-body problem
 - Multiple scattering expansion $d+A$ to be developed

(d,p) Reaction as
three-body problem:

Faddeev AGS



Elastic, breakup, rearrangement channels
are included and fully coupled.

Deltuva and Fonseca, Phys. Rev. **C79**, 014606 (2009).

Hamiltonian: $\mathbf{H} = \mathbf{H}_0 + \mathbf{V}_{np} + \mathbf{V}_{nA} + \mathbf{V}_{pA}$

Issues:

- Traditional Faddeev eqs do not consider target excitations
- Coulomb force for heavy nuclei: current calculations do not converge for $Z \geq 20$

Suggestion to extend Faddeev equations:

- Inclusion of target excitations
- Novel treatment of Coulomb in momentum space

A.M. Mukhamedzhanov, V.Eremenko and A.I. Sattarov,
arXiv:1206.3791 [nucl-th]

Including target excitations:

Formulation natural when transition amplitudes in sub-systems are separable.

Reminder:

t-matrices (= interactions summed up to all orders) from two-body subsystems are input to Faddeev amplitudes

Optical Potentials (phenomenological) for proton and neutron elastic scattering from nuclei:

Wood-Saxon functions in coordinate space

Example: Central part of CH89 optical potential

$$U_{\text{nucl}}(r) = V(r) + i(W(r) + W_s(r)).$$

$$V(r) = \frac{-V_r}{1 + \exp\left(\frac{r-R_r}{a_r}\right)}$$

$$W(r) = \frac{-V_i}{1 + \exp\left(\frac{r-R_i}{a_i}\right)}$$

$$W_s(r) = \frac{-V_s \times 4 \times \exp\left(\frac{-(r-R_s)}{a_s}\right)}{\left(1 + \exp\left(\frac{-(r-R_s)}{a_s}\right)\right)^2}$$

Not a form that can be used in a Faddeev calculation in momentum space

Our test case:

$n+^{48}\text{Ca}$ CH89 central optical potential

Steps for creating a separable t-matrix:

1. Fourier Transform

Analytic in form of a series

– lowest terms sufficient.

E_{lab} [MeV]	δ_0 [deg]	
	k – space	r – space
5	12.645	12.644
10	-41.514	-41.516
20	-74.221	-74.220
40	-0.130	-0.132
50	-26.0297	-26.0296

2. Separabilization via EST scheme

D. J. Ernst, C. M. Shakin and R. M. Thaler, Phys. Rev. C 8, 46 (1973).

Use in momentum space

Use half-shell t-matrices as form factors at EST support points

EST Scheme:

$$\hat{t}(p, p', E) = \sum_{ij} \langle p | v_l | \psi_i \rangle \tau_{ij}(E) \langle \psi_j | v_l | p' \rangle$$

i, j give the rank

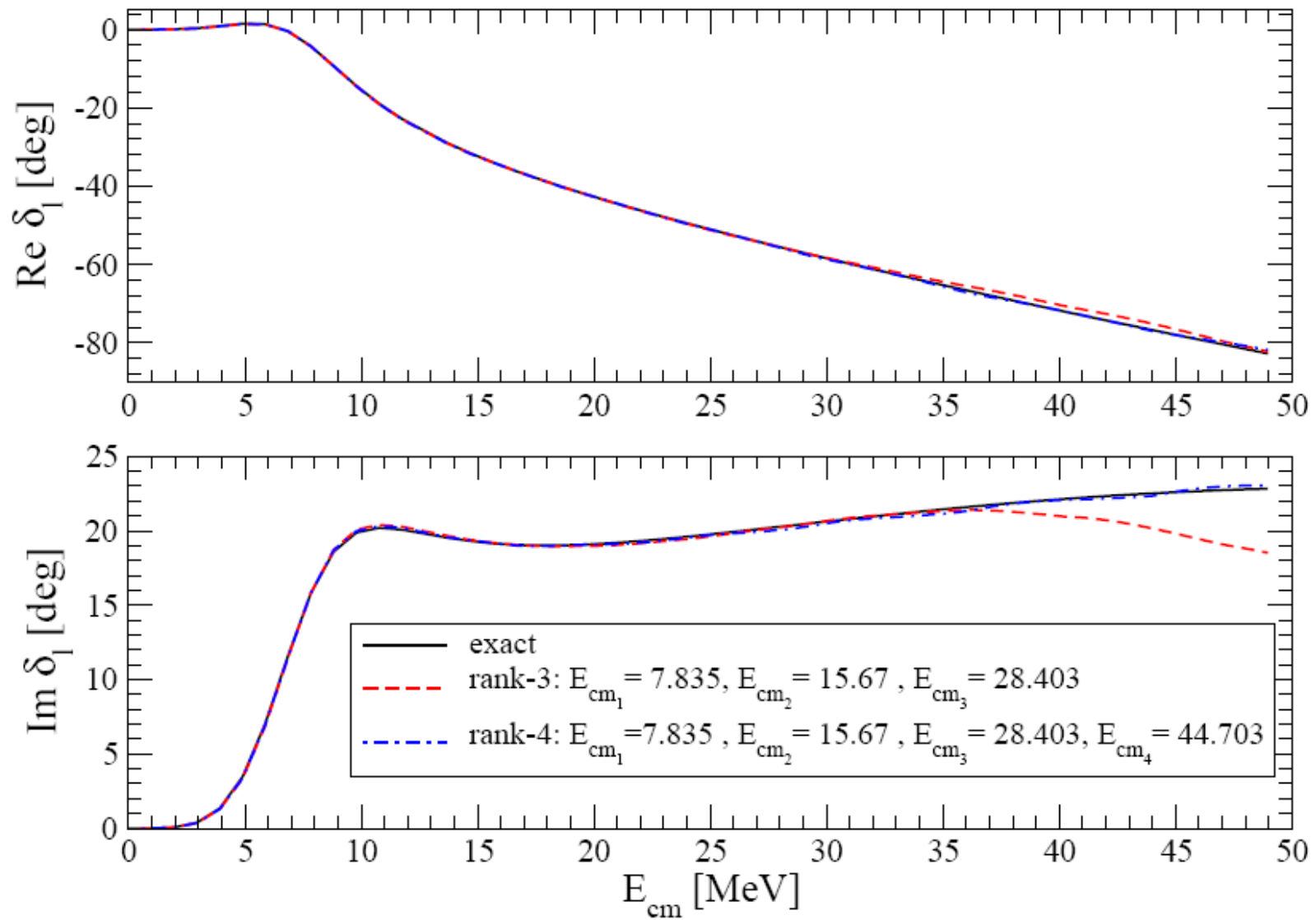
$$\sum_j \tau_{ij}(E) \underbrace{\langle \Psi_j | v - v g_0(E) v | \Psi_k \rangle}_{\text{matrix element}} = \delta_{ik}$$

*Calculate matrix for given j, k
and solve algebraic system of eqs.*

- $v | \Psi_j \rangle :=$ half-shell t-matrices at fixed energies E_k
- t-matrix exact those $E_k =:$ EST support points
- EST separable t-matrix interpolates between the support points

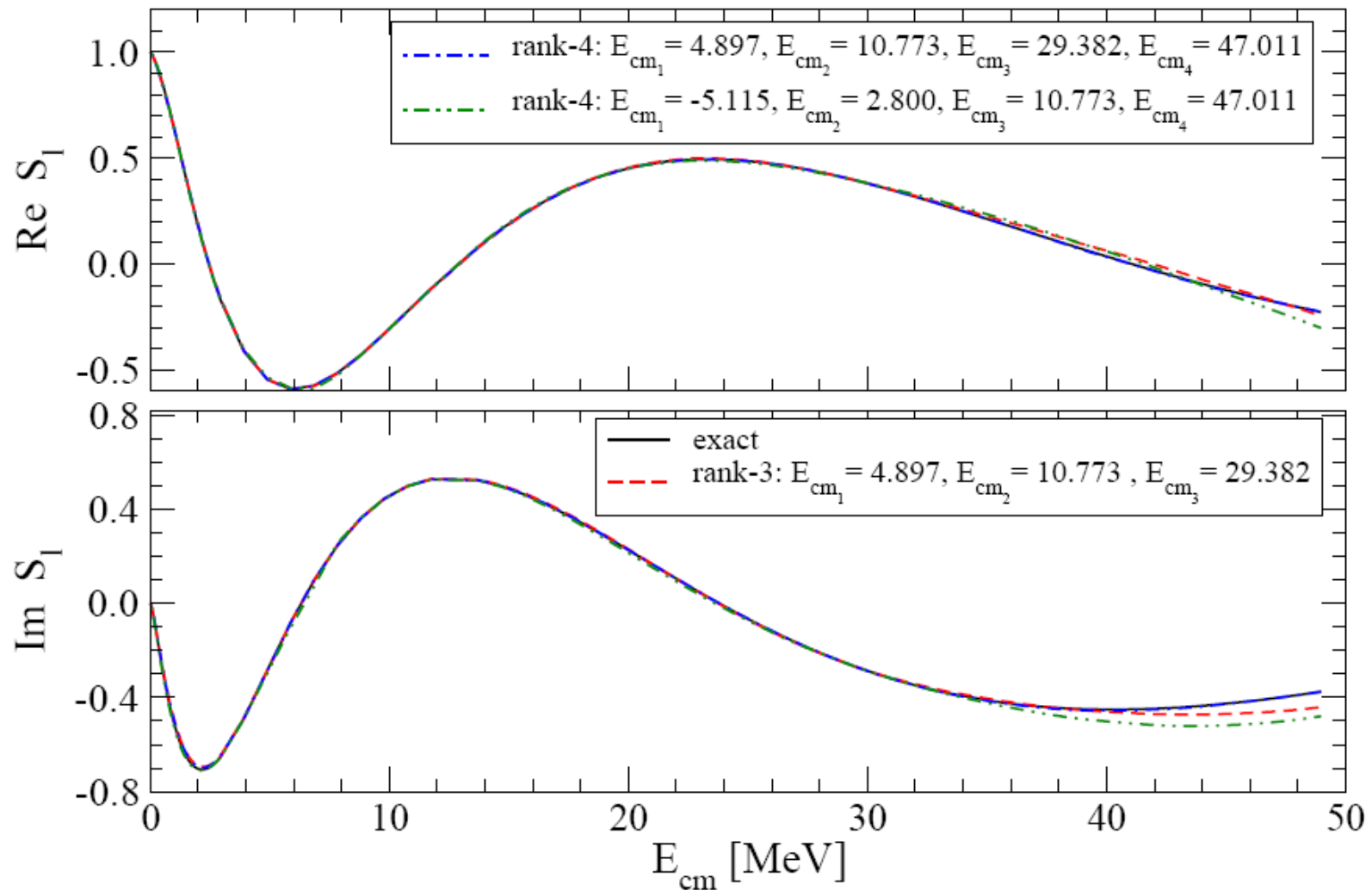
$n\text{-}^{48}\text{Ca}$ phase shifts : $l = 4$ partial wave

CH89 Optical Potential - Central



$n\text{-}^{48}\text{Ca}$ s-matrix : $l = 1$ partial wave

CH89 Optical Potential - Central



Next: spin-orbit ... straightforward ...